

# **Orthogonal Functions & Inner Product**

- □ Vectors in linear algebra are not just *n*-tuples
- □ If u and v are two *n*-tuple vectors in 3D-space, then the inner product (u, v) possesses the following properties:
  - $(\mathbf{u}, \mathbf{v}) = (\mathbf{v}, \mathbf{u})$  (inner product is commutable)
  - $(k\mathbf{u}, \mathbf{v}) = k(\mathbf{u}, \mathbf{v})$  (k is a scalar)
  - $(\mathbf{u}, \mathbf{u}) = 0$ , if  $\mathbf{u} = 0$  and  $(\mathbf{u}, \mathbf{u}) > 0$ , if  $\mathbf{u} \neq 0$
  - $(\mathbf{u}+\mathbf{v}, \mathbf{w}) = (\mathbf{u}, \mathbf{w}) + (\mathbf{v}, \mathbf{w})$  (inner product is distributable)
- □ The definite integral of two functions,  $f_1$  and  $f_2$ , over an interval [a, b] possesses the same properties as well → we can define "inner product" for functions

## **Inner Product of Functions**

□ The inner product of two functions  $f_1$  and  $f_2$  on an interval [a, b] is the number

$$(f_1, f_2) = \int_a^b f_1(x) f_2(x) dx$$

□ Two vectors are "orthogonal" if the inner product is zero
 → function inner product should be defined similarly:
 two functions f<sub>1</sub> and f<sub>2</sub> are orthogonal on an interval
 [a, b] if

$$(f_1, f_2) = \int_a^b f_1(x) f_2(x) dx = 0$$

# **Orthonormal Set of Functions**

□ A set of real-valued functions { $\phi_0$ ,  $\phi_1$ ,  $\phi_2$ , ...} is said to be orthogonal on an interval [a, b] if

$$(\phi_m,\phi_n) = \int_a^b \phi_m(x)\phi_n(x)dx = 0, \quad m \neq n.$$

□ The norm of a function  $\phi$  is defined as  $||\phi|| = (\phi, \phi)^{1/2}$ . That is,

$$\left|\phi(x)\right\| = \sqrt{\int_a^b \phi^2(x) dx}.$$

If  $\{\phi_n\}$  is an orthogonal set on [a, b] and  $\|\phi_n\| = 1$ ,  $\forall n$ , then  $\{\phi_n\}$  is an orthonormal set of functions on [a, b].

## **Orthogonal Series Expansion**

□ If  $\{\phi_n(x)\}$  is an infinite orthogonal set of functions on the interval [a, b], is it possible to determine a set of coefficients  $c_n$ , n = 0, 1, 2, ... such that

 $f(x) = c_0 \phi_0(x) + c_1 \phi_1(x) + \dots + c_n \phi_n(x) + \dots ?$ 

To find the coefficient of  $\phi_n$ , we compute  $(f, \phi_n)$ 

 $(f, \phi_n) = c_0(\phi_0, \phi_n) + c_1(\phi_1, \phi_n) + \ldots + c_n(\phi_n, \phi_n) + \ldots$ Since  $\{\phi_n(x)\}$  is an orthogonal set,  $(\phi_m, \phi_n) = 0, \forall m \neq n$ . Therefore,

$$c_n = \frac{(f, \phi_n)}{\|\phi_n\|^2}$$
 and  $f(x) = \sum_{n=0}^{\infty} \frac{(f, \phi_n)}{\|\phi_n\|^2} \phi_n(x).$ 

# **Completeness of an Orthogonal Set**

□ In previous discussion, we have  $f(x) = \sum_{n=0}^{\infty} \frac{(f, \phi_n)}{\|\phi_n\|^2} \phi_n(x)$ , if f(x) can be represented as a linear combination of  $\phi_0(x) \sim \phi_\infty(x)$  in the vector space *S*.

However, not every functions in *S* can be represented as a linear combinations of the functions in  $\{\phi_n(x)\}$ . This is true only when  $\{\phi_n(x)\}$  is a complete set of *S*, i.e., when  $\{\phi_n(x)\}$  is a vector basis of *S*.

## **Periodic External Forces**

□ Recall that a linear 2<sup>nd</sup>-order DE:

$$\frac{d^2x}{dt^2} + \omega_0^2 x = f(t),$$

where f(t) stands for the external force imposed on the (undamped) system. Often, f(t) is a periodic function (over an interval of interest).

- Question: Is there a systematic way to represent a general periodic function?
  - Well, Taylor series may work, but can we do better?

# Properties of a Periodic Function

- □ **Definition:** The function f(t) defined for all *t* is said to be periodic provided that there exists a positive number *p* such that f(t + p) = f(t) for all *t*. If *p* is the smallest number with this property, then *p* is called the period of the function *f*.
- □ Remarks:
  - Linear combinations of two (or more) periodic functions will still be a periodic function.
  - If we use a set of periodic functions as basis functions to represent other periodic functions, they should work better than if we use {1, x, x<sup>2</sup>, x<sup>3</sup>, ...}, as in Taylor series.

# **Selection of Periodic Basis**

□ In 1822, J. Fourier asserted that every function f(t) with period  $2\pi$  can be represented as a linear combination of sin *nt* and cos *nt*, as follows:

$$\frac{a_0}{2} + \sum_{n=1}^{\infty} (a_n \cos nt + b_n \sin nt),$$

□ Really? How about the function:



### **Fourier Series**

□ Note that the set of trigonometric functions:

{1, cos *t*, cos 2*t*, cos 3*t*, ..., sin *t*, sin 2*t*, sin3*t*, ...}
are orthogonal on the interval [-*π*, *π*].
□ The Fourier series of *f*(*t*) on [-*π*, *π*] is defined as :

where 
$$f(t) = \frac{a_0}{2} + \sum_{n=1}^{\infty} (a_n \cos nt + b_n \sin nt),$$
  
The projection vector of  $f(t)$  onto  $\cos nt$  is  
 $a_0 = \frac{1}{\pi} \int_{-\pi}^{\pi} f(t) dt$ 
 $\frac{\langle f(t), \cos nt \rangle}{\|\cos nt\|^2} \cos nt.$  Thus,  $a_n = \frac{\langle f(t), \cos nt \rangle}{\|\cos nt\|^2}$   
 $a_n = \frac{1}{\pi} \int_{-\pi}^{\pi} f(t) \cos nt dt,$ 
 $b_n = \frac{1}{\pi} \int_{-\pi}^{\pi} f(t) \sin nt dt.$ 



Example: 
$$f(x) = \begin{cases} 0, & -\pi < x < 0 \\ \pi - x, & 0 \le x < \pi \end{cases}$$
 (1/2)

 $\Box$  Since  $p = \pi$ , we have

$$a_{0} = \frac{1}{\pi} \int_{-\pi}^{\pi} f(x) dx$$

$$= \frac{1}{\pi} \left[ \int_{-\pi}^{0} 0 dx + \int_{0}^{\pi} (\pi - x) dx \right]$$

$$= \frac{1}{\pi} \left[ \pi x - \frac{x^{2}}{2} \right]_{0}^{\pi} = \frac{\pi}{2}$$

Example: 
$$f(x) = \begin{cases} 0, & -\pi < x < 0 \\ \pi - x, & 0 \le x < \pi \end{cases}$$
 (2/2)  
$$a_n = \frac{1}{\pi} \int_{-\pi}^{\pi} f(x) \cos nx dx = \frac{1}{\pi} \left[ \int_{-\pi}^{0} 0 dx + \int_{0}^{\pi} (\pi - x) \cos nx dx \right]$$
$$= \frac{1}{\pi} \left[ (\pi - x) \frac{\sin nx}{n} \right]_{0}^{\pi} + \frac{1}{n} \int_{0}^{\pi} \sin nx dx \right]$$
$$= \frac{1}{\pi} \frac{\cos nx}{n} \int_{0}^{\pi} \frac{-\cos n\pi + 1}{n^{2}\pi} \qquad (\text{Note that } \cos n\pi = (-1)^{n})$$
$$= \frac{1}{\pi} \int_{0}^{\pi} (\pi - x) \sin nx dx = \frac{1}{n}$$
$$f(x) = \frac{\pi}{4} + \sum_{n=1}^{\infty} \left\{ \frac{1 - (-1)^{n}}{n^{2}\pi} \cos nx + \frac{1}{n} \sin nx \right\}$$

# Fourier Convergence Theorem

□ Theorem: Let f and f' be piecewise continuous on the interval (-p, p); that is f and f' be continuous except at a finite number of points, then the Fourier series of f converges to f at a point of continuity.

At a point of discontinuity the Fourier series converges to the average:

$$\frac{f(x+)+f(x-)}{2},$$

where f(x+) and f(x-) denote the limit of f at x from the right and the left, respectively

## Example: Converges at Discontinuity

**\Box** The following function is discontinuous at x = 0:

$$f(x) = \begin{cases} 0, & -\pi < x < 0\\ \pi - x, & 0 \le x < \pi \end{cases}$$

□ The series converges to f at  $x \neq 0$ . At x = 0, the series converges to:

$$\frac{f(0+)+f(0-)}{2} = \frac{\pi+0}{2} = \frac{\pi}{2}$$

## **Periodic Extension**

- □ Fourier series not only represents a function *f* on the interval (−*p*, *p*), but also gives the periodic extension of *f* outside the interval.
- □ When *f* is piecewise continuous and the right- and lefthand derivatives exist at x = -p and x = p, respectively, then the series converges to the average [f(-p-) + f(-p+)]/2 = [f(p-) + f(-p+)]/2 at the end points:









# Properties of Even/Odd Functions

- □ The product of two even functions is even
- □ The product of two odd functions is even
- □ The product of an even and an odd functions is odd
- □ The sum (difference) of two even functions is even
- The sum (difference) of two odd functions is odd
  If *f* is even, then

$$\int_{-a}^{a} f(t)dt = 2\int_{0}^{a} f(t)dt.$$

 $\Box$  If *f* is odd, then

$$\int_{-a}^{a} f(t)dt = 0.$$

## **Cosine and Sine Series**

 $\Box$  If *f* is an even function on (–*p*, *p*), then

$$a_{0} = \frac{1}{p} \int_{-p}^{p} f(x) dx = \frac{2}{p} \int_{0}^{p} f(x) dx,$$
  

$$a_{n} = \frac{1}{p} \int_{-p}^{p} f(x) \cos \frac{n\pi}{p} x dx = \frac{2}{p} \int_{0}^{p} f(x) \cos \frac{n\pi}{p} x dx,$$
  

$$b_{n} = \frac{1}{p} \int_{-p}^{p} f(x) \sin \frac{n\pi}{p} x dx = 0.$$

 $\Box$  Similarly, if *f* is odd,

$$a_n = 0, n = 0, 1, 2, ..., \quad b_n = \frac{2}{p} \int_0^p f(x) \sin \frac{n\pi}{p} x dx.$$

## Fourier Cosine and Sine Series

□ Suppose that the function f(x) is piecewise continuous on the interval [0, p]. The Fourier cosine series of *f* is:

$$f(x) = \frac{a_0}{2} + \sum_{n=1}^{\infty} a_n \cos \frac{n\pi}{p} x$$
, with  $a_n = \frac{2}{p} \int_0^p f(x) \cos \frac{n\pi}{p} x dx$ .

The Fourier sine series of f is:

$$f(x) = \sum_{n=1}^{\infty} b_n \sin \frac{n\pi}{p} x, \text{ with } b_n = \frac{2}{p} \int_0^p f(x) \sin \frac{n\pi}{p} x dt.$$

Example: 
$$f(x) = \begin{cases} -1, & -\pi < x < 0 \\ 0, & x = 0, \pi, -\pi \\ +1, & 0 < x < \pi \end{cases}$$

 $\Box$  Calculate  $b_n$  as follows:



# Half-Range Expansions

□ Sometimes, we only care about the Fourier series defined on (0, L). We can define the function f on (-L, 0) so that the expansion has a simpler form.

□ Three possible choices of extension:



The third one has period *L*, others have period 2*L*.

Example: 
$$f(x) = 2x - x^2, x \in (0, 2)$$

□ We can expand f(x) to the range (-2, 2) and make it an even ( $f_{even}(x)$ ) or an odd ( $f_{odd}(x)$ ) function:

$$f_{\text{even}}(x) = f(-x) = 2(-x) - (-x)^2 = -2x - x^2$$
, for  $x < 0$ ,

or

$$f_{\text{odd}}(x) = -f(-x) = -[2(-x) - (-x)^2] = 2x + x^2$$
, for  $x < 0$ .



The Fourier expansion of  $f_{even}(x)$  has only cosine terms while  $f_{odd}(x)$  has only sine terms.

**Example:** 
$$f(x) = x^2, 0 < x < L$$

 $\Box$  Expand f(x) in a (a) cosine, (b) sine, (c) Fourier series



# **Review: Periodic Driving Force**

When the driving force f(t) of a DE is periodic and defined over [0, p], Half-range expansion of Fourier series are quite useful. For example, the particular solution of the DE:

$$m\frac{d^2x}{dt^2} + kx = f(t)$$

can be solved by first representing f(t) by a half-range sine expansion and assume a particular solution of the form:

$$x_p(t) = \sum_{n=1}^{\infty} B_n \sin \frac{n\pi}{p} t$$

Example: 
$$x''+4x = 4t$$
,  $x(0)=x(1)=0$  (1/2)

□ Assume that 0 < t < 1 for f(t), we can use odd extension with p = 1 to get the Fourier sine series of f:

$$4t = \frac{8}{\pi} \sum_{n=1}^{\infty} \frac{(-1)^{n+1}}{n} \sin n \pi t.$$

The solution x(t) should be in sine series form as well:

$$x(t) = \sum_{n=1}^{\infty} b_n \sin n \pi t.$$

Note that x(t) satisfies the boundary conditions. Substitute the solution into the DE, we have

$$\sum_{n=1}^{\infty} (-n^2 \pi^2 + 4) b_n \sin n \pi t = \frac{8}{\pi} \sum_{n=1}^{\infty} \frac{(-1)^{n+1}}{n} \sin n \pi t.$$

Example: 
$$x''+4x = 4t$$
,  $x(0)=x(1)=0$  (2/2)

□ The solution of the coefficients  $b_n$  is then

$$b_n = \frac{8 \cdot (-1)^{n+1}}{n \pi (4 - n^2 \pi^2)}.$$

The Fourier series solution can be expressed as:

$$x(t) = \frac{8}{\pi} \sum_{n=1}^{\infty} \frac{(-1)^{n+1} \sin n \pi t}{n(4 - n^2 \pi^2)}, \quad (0 \le t \le 1)$$

which is equivalent to

$$x(t) = t - \frac{\sin 2t}{\sin 2}$$

in the interval (-1, 1).

