

Power Series

 \Box A power series in (x - a) is a series of the form

$$\sum_{n=0}^{\infty} c_n (x-a)^n = c_0 + c_1 (x-a) + c_2 (x-a)^2 + \cdots.$$

Such a series is said to be a power series centered at *a*.

□ A power series is convergent at a value $x \in I$ if the limit of partial sums exists, i.e.

$$f(x) = \lim_{N \to \infty} \sum_{n=0}^{N} c_n (x-a)^n$$

The interval of convergence, *I*, of a power series is the set of all numbers where the series converges.

Ratio Test

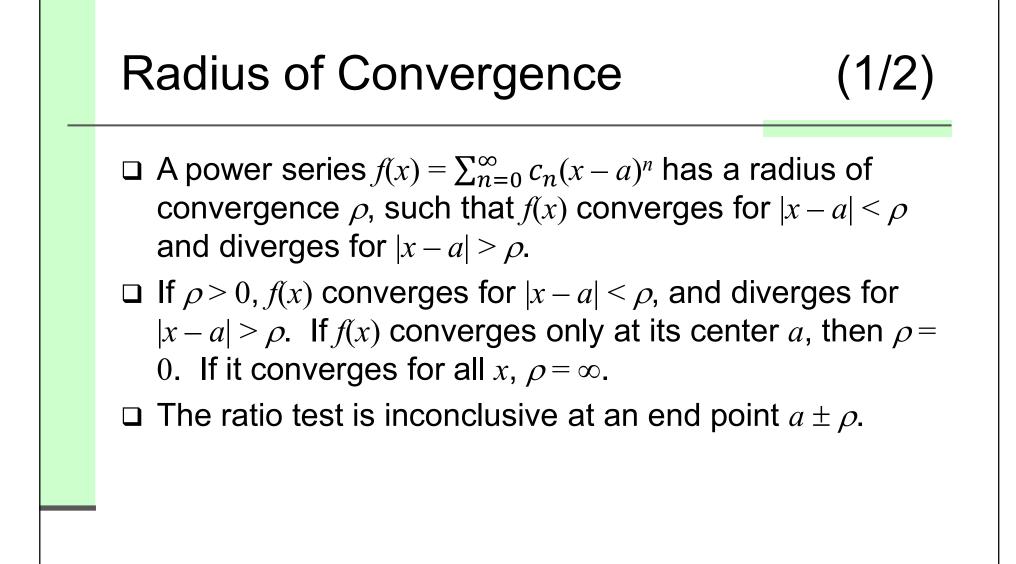
□ Convergence of a power series can often be checked by the ratio test: suppose $c_n \neq 0$ for all *n* in

$$\sum_{n=0}^{\infty} c_n (x-a)^n,$$

and that

$$\lim_{n \to \infty} \left| \frac{c_{n+1} (x-a)^{n+1}}{c_n (x-a)^n} \right| = |x-a| \lim_{n \to \infty} \left| \frac{c_{n+1}}{c_n} \right| = L.$$

If L < 1, the series converges absolutely; if L > 1 the series diverges; and if L = 1 the test is inconclusive.



Radius of Convergence

(2/2)

 \Box Given the power series $g(x) = \sum_{n=0}^{\infty} c_n x^n$, if the limit

$$\rho = \lim_{n \to \infty} \left| \frac{c_n}{c_{n+1}} \right|$$

exists, then

- If $\rho = 0$, then g(x) diverges for all $x \neq 0$.
- If $0 < \rho < \infty$, g(x) converges if $|x| < \rho$, and diverges if $|x| > \rho$.
- If $\rho = \infty$, then g(x) converges for all x.

Power Series of a Function f(x)

 \Box The Taylor series of a function f(x) is defined as

$$y = \sum_{n=0}^{\infty} \frac{f^{(n)}(a)}{n!} (x-a)^n.$$

If *y* converges to f(x) for all *x* in some open interval containing *a*, then we say that the function f(x) is analytic at x = a.

 Polynomials are analytic; a rational function is analytic wherever the denominator is not zero.

□ Arithmetic of power series

- The operations of addition, multiplication, and division can be applied to power series as in polynomials.
- If *f* and *g* are analytic at *a*, so are *f*+g, *f*·*g*, and *f*/*g* (if $g(a) \neq 0$).

Examples of Power Series

□ By Taylor ($\forall a$) or Maclaurin (a = 0) series expansions, common functions can be written in power series forms:

$$e^{x} = \sum_{n=0}^{\infty} \frac{x^{n}}{n!} = 1 + x + \frac{x^{2}}{2!} + \frac{x^{3}}{3!} + \cdots$$

$$\sin x = \sum_{n=0}^{\infty} \frac{(-1)^{n} x^{2n+1}}{(2n+1)!} = x - \frac{x^{3}}{3!} + \frac{x^{5}}{5!} - \cdots$$

$$\ln(1+x) = \sum_{n=1}^{\infty} \frac{(-1)^{n+1} x^{n}}{n!} = x - \frac{x^{2}}{2} + \frac{x^{3}}{3} - \cdots$$

$$\frac{1}{1-x} = \sum_{n=0}^{\infty} x^{n} = 1 + x + x^{2} + x^{3} + \cdots$$

• •

Example: Adding Two Power Series

□ Write $\sum_{n=2}^{\infty} n(n-1)c_n x^{n-2} + \sum_{n=0}^{\infty} c_n x^{n+1}$ as one power series.

Solution:

$$\begin{split} &\sum_{n=2}^{\infty} n(n-1)c_n x^{n-2} + \sum_{n=0}^{\infty} c_n x^{n+1} = 2 \cdot 1 \cdot c_2 x^0 + \sum_{n=3}^{\infty} n(n-1)c_n x^{n-2} + \sum_{n=0}^{\infty} c_n x^{n+1} \\ &= 2c_2 + \sum_{k=1}^{\infty} (k+2)(k+1)c_{k+2} x^k + \sum_{k=1}^{\infty} c_{k-1} x^k \\ &= 2c_2 + \sum_{k=1}^{\infty} [(k+2)(k+1)c_{k+2} + c_{k-1}]x^k. \end{split}$$

Power Series Method

The power series method for solving a DE consists of substituting the power series

$$y = \sum_{n=0}^{\infty} c_n x^n$$

in the DE and determining the coefficients c_0, c_1, \ldots so that the equation satisfies.

$$\Box \text{ If } f(x) = \sum_{n=0}^{\infty} c_n x^n \text{, then } f'(x) = \sum_{n=1}^{\infty} n c_n x^{n-1}$$

□ If $\sum_{n=0}^{\infty} a_n x^n = \sum_{n=0}^{\infty} b_n x^n$, for all *x* in the interval of convergence, then $a_n = b_n$ for all $n \ge 0$.

Example:
$$y' + 2y = 0$$

$$\Box$$
 Since $y = \sum_{n=0}^{\infty} c_n x^n$, and $y' = \sum_{n=1}^{\infty} n c_n x^{n-1}$

we have

$$\sum_{n=1}^{\infty} nc_n x^{n-1} + 2\sum_{n=0}^{\infty} c_n x^n = 0.$$

Perform change of index *n* to align x^n , $\sum_{n=0}^{\infty} (n+1)c_{n+1}x^n + 2\sum_{n=0}^{\infty} c_n x^n = 0 \rightarrow \sum_{n=0}^{\infty} [(n+1)c_{n+1} + 2c_n]x^n = 0.$

We have a recurrence relation $c_{n+1} = -2c_n/(n+1)$, $n \ge 0$.

$$\rightarrow c_n = (-2)^n c_0 / n!, \ n \ge 1 \ \rightarrow \ y(x) = \sum_{n=0}^{\infty} \frac{(-2)^n c_0}{n!} x^n.$$

Power Series Solutions

□ Suppose the linear 2nd-order DE

 $a_2(x)y'' + a_1(x)y' + a_0(x)y = 0$

is put into the standard form

$$y'' + P(x)y' + Q(x)y = 0,$$

then:

A point $x = x_0$ is said to be an ordinary point of the DE if both P(x) and Q(x) are analytic at x_0 . A point that is not an ordinary point is said to be a singular point of the equation.

Example: Ordinary, Singular Points

 \Box Every finite value of *x* is an ordinary point of

 $y'' + (e^x) y' + (\sin x) y = 0.$

 \Box *x* = 0 is a singular point of the DE

 $y'' + (e^x) y' + (\ln x) y = 0.$

Polynomial-Coefficient DEs

□ Recall that a polynomial is analytic at any value *x*, and a rational function is analytic except at points where its denominator is zero. Thus, a 2nd-order polynomialcoefficient DE has singular points when $a_2(x) = 0$, since $a_1(x) = a_2(x)$

$$P(x) = \frac{a_1(x)}{a_2(x)}, \quad Q(x) = \frac{a_0(x)}{a_2(x)}.$$

□ Examples

- The Euler equation $ax^2y'' + bxy' + cy = 0$ has a singular point at x = 0.
- The equation $(x^2 + 1)y'' + xy' y = 0$ has singular points at $x = \pm i$.

Solutions Near an Ordinary Point

Theorem: If $x = x_0$ is an ordinary point of

 $a_2(x)y'' + a_1(x)y' + a_0(x)y = 0,$

we can always find two linearly independent solutions in the form of a power series centered at x_0 , that is,

$$y = \sum_{n=0}^{\infty} c_n (x - x_0)^n.$$

A series solution converges at least on some interval defined by $|x - x_0| < \rho$, where ρ is the distance from x_0 to the nearest singular point.

□ Note that for $|x - x_0| \ge \rho$, y(x) may or may not converge. Further investigations are required.

Ex:
$$(x^2 - 4)y'' + 3xy' + y = 0, y(0) = 4, y'(0) = 1$$

Note that the singular points are ±2, there should be a solution with radius of convergence at least 2. Since

$$y = \sum_{n=0}^{\infty} c_n x^n$$
, $y' = \sum_{n=1}^{\infty} n c_n x^{n-1}$, and $y'' = \sum_{n=2}^{\infty} n(n-1) c_n x^{n-2}$

Therefore,

$$\sum_{n=2}^{\infty} n(n-1)c_n x^n - 4\sum_{n=2}^{\infty} n(n-1)c_n x^{n-2} + 3\sum_{n=1}^{\infty} nc_n x^n + \sum_{n=0}^{\infty} c_n x^n = 0.$$

$$\sum_{n=0}^{\infty} n(n-1)c_n x^n - 4\sum_{n=0}^{\infty} (n+2)(n+1)c_{n+2} x^n + 3\sum_{n=0}^{\infty} nc_n x^n + \sum_{n=0}^{\infty} c_n x^n = 0.$$

Combining the terms, we have
$$c_{n+2} = \frac{(n+1)c_n}{4(n+2)}$$
, $n \ge 0$.

Ex:
$$(x^2 - 4)y'' + 3xy' + y = 0$$
 (2/2)

When $n = 0, 2, 4, ...,$ we have
 $c_2 = \frac{c_0}{4 \cdot 2}, c_4 = \frac{3c_0}{4^2 \cdot 2 \cdot 4}, c_6 = \frac{3 \cdot 5c_0}{4^3 \cdot 2 \cdot 4 \cdot 6}, \dots \rightarrow c_{2n} = \frac{1 \cdot 3 \cdot 5 \cdots (2n-1)}{4^n \cdot 2 \cdot 4 \cdots (2n)} c_0.$
When $n = 1, 3, 5, ...,$ we have
 $c_3 = \frac{2c_1}{4 \cdot 3}, c_5 = \frac{2 \cdot 4c_1}{4^2 \cdot 3 \cdot 5}, c_7 = \frac{2 \cdot 4 \cdot 6c_1}{4^3 \cdot 3 \cdot 5 \cdot 7}, \dots \rightarrow c_{2n+1} = \frac{2 \cdot 4 \cdot 6 \cdots (2n)}{4^n \cdot 1 \cdot 3 \cdots (2n+1)} c_1.$
Therefore, the solution is
 $y(x) = c_0 \left(1 + \sum_{n=1}^{\infty} \frac{1 \cdot 3 \cdot 5 \cdots (2n-1)}{2^{3n} n!} x^{2n} \right) + c_1 \left(x + \sum_{n=1}^{\infty} \frac{n!}{2^n \cdot 1 \cdot 3 \cdots (2n+1)} x^{2n+1} \right).$
and $y(0) = c_0 = 4, y'(0) = c_1 = 1.$

Translated Series Solutions

□ If the initial condition is given at x_0 other than zero, we have to assume the general solution form

$$y(x) = \sum_{n=0}^{\infty} c_n (x - x_0)^n.$$

This way, we can obtain the IVP solution with $y(x_0) = c_0$ and $y'(x_0) = c_1$ easily.

□ As an alternative, we can translate the equation by letting $t = x - x_0$ and assume the solution form:

$$y = \sum_{n=0}^{\infty} c_n x^n$$

EX:
$$(t^2-2t-3)y''+3(t-1)y'+y=0, y(1)=4, y'(1)=-1$$

□ Perform a change of variable to the DE by x = t - 1:

$$t^2 - 2t - 3 = (x + 1)^2 - 2(x + 1) - 3 = x^2 - 4$$
,

and

$$\frac{dy}{dt} = \frac{dy}{dx}\frac{dx}{dt} = \frac{dy}{dx}, \quad \frac{d^2y}{dt^2} = \left[\frac{d}{dx}\left(\frac{dy}{dx}\right)\right]\frac{dx}{dt} = \frac{d^2y}{dx^2}.$$

Hence, the DE becomes $(x^2-4)d^2y/dx^2 + 3x(dy/dx) + y = 0$. \rightarrow Same DE as the previous one.

Substituting x = t-1 into the previous solution, we get

$$y(t) = 4 + (t-1) + \frac{1}{2}(t-1)^2 + \frac{1}{6}(t-1)^3 + \frac{3}{32}(t-1)^4 + \cdots$$

which converges if -1 < t < 3.

Example: $y'' + (\cos x) y = 0$ (1/2)

□ Since

 $y'' + (\cos x)y$ $= \sum_{n=2}^{\infty} n(n-1)c_n x^{n-2} + \left(1 - \frac{x^2}{2!} + \frac{x^4}{4!} - \frac{x^6}{6!} + \cdots\right) \sum_{n=0}^{\infty} c_n x^n$ $= 2c_2 + 6c_3 x + 12c_4 x^2 + \cdots + \left(1 - \frac{x^2}{2!} + \frac{x^4}{4!} + \cdots\right) \left(c_0 + c_1 x + c_2 x^2 + \cdots\right)$ $= 2c_2 + c_0 + (6c_3 + c_1)x + \left(12c_4 + c_2 - \frac{1}{2}c_0\right)x^2 + \cdots = 0$

Example:
$$y'' + (\cos x) y = 0$$
 (2/2)

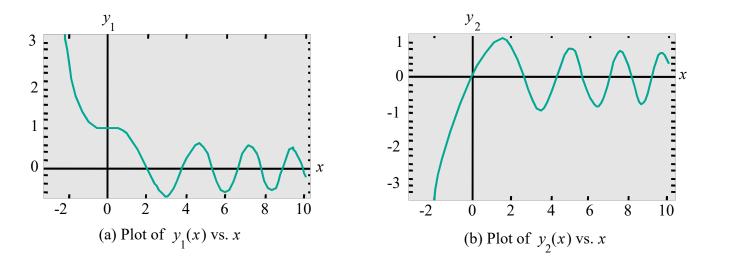
We have:

$$2c_2 + c_0 = 0$$
, $6c_3 + c_1 = 0$, $12c_4 + c_2 - \frac{1}{2}c_0 = 0$, ...

Therefore:

$$y_1(x) = 1 - \frac{1}{2}x^2 + \frac{1}{12}x^4 - \dots, \ y_2(x) = x - \frac{1}{6}x^3 + \frac{1}{30}x^5 - \dots$$

with region of convergence $|x| < \infty$.



Solutions about Singular Points

□ Let $y(x) = \Sigma c_n x^n$, if P(x) is not analytic at 0, its power series form $\Sigma b_k x^k$ will not converge to P(0) at 0 given any b_k . However, it is possible that

 $(\Sigma b_k x^k) (\Sigma n c_n x^{n-1})$

may still converge to P(x)y'(x).

In short, even if x = 0 is a singular point, the power series expression of

y'' + P(x)y' + Q(x)y

may still converge to zero.

Regular Singular Points

□ Assume that a DE in the standard form

y'' + P(x)y' + Q(x)y = 0 has a singular point at x_0 .

If there are two functions p(x) and q(x), both are analytic at x_0 , such that $p(x) = (x - x_0)P(x)$ and $q(x) = (x - x_0)^2Q(x)$, the original DE can be rewritten as:

$$y'' + \frac{p(x)}{(x - x_0)}y' + \frac{q(x)}{(x - x_0)^2}y = 0,$$

then, we call $x = x_0$ a regular singular point of the DE.

Otherwise, $x = x_0$ is a irregular singular point.

Remarks on Singularity of *P* and *Q*

- □ If x = 0 is a singular point, the power series expansion of P(x) at 0 approaches ∞.
- □ However, if P(x) grows slower than 1/x when $x \to 0$, then xP(x) is convergent. That is, p(x) = xP(x) is analytic at 0. Similarly, q(x) is analytic at 0 if Q(x) grows slower than $1/x^2$.

□ Note that, for the DE
$$y'' + \frac{p(x)}{x}y' + \frac{q(x)}{x^2}y = 0$$
,

x = 0 is a regular singular point if p(x) and q(x) are polynomials.

Example: Singular Points

□ For $(x^2 - 4)^2 y'' + 3(x - 2)y' + 5y = 0$, x = 2 and x = -2 are singular points. We have

$$P(x) = \frac{3}{(x-2)(x+2)^2} \quad \text{and} \quad Q(x) = \frac{5}{(x-2)^2(x+2)^2}.$$

□ Obviously x = 2 is a regular singular point, and x = -2 is an irregular singular point.

Example: Non-polynomial p(x), q(x)

 $\Box \text{ The DE } x^4 y'' + (x^2 \sin x)y' + (1 - \cos x)y = 0 \text{ can be}$ expressed as $y'' + \frac{\sin x/x}{x}y' + \frac{(1 - \cos x)/x^2}{x^2}y = 0.$ Since x = 0 is not an ordinary point and $p(x) = \frac{\sin x}{x} = \frac{1}{x} \left(x - \frac{x^3}{3!} + \frac{x^5}{5!} - \cdots \right) = 1 - \frac{x^2}{3!} + \frac{x^4}{5!} - \cdots,$ $q(x) = \frac{1 - \cos x}{x^2} = \frac{1}{x^2} \left[1 - \left(1 - \frac{x^2}{2!} + \frac{x^4}{4!} - \cdots \right) \right] = \frac{1}{2!} - \frac{x^2}{4!} + \frac{x^4}{6!} \cdots,$

are both analytic (convergent) at 0, thus x = 0 is a regular singular point.

Solution near Singular Points

□ For a constant-coefficient Cauchy-Euler equation $x^2y'' + p_0xy' + q_0y = 0$,

where p_0 and q_0 are constants, we can assume that $y(x) = x^r$ is a solution $\rightarrow r$ is a root of the equation: $r(r-1) + p_0r + q_0 = 0.$

If we have coefficient functions p(x) and q(x) instead, is it possible that

$$y(x) = x^r \sum_{n=0}^{\infty} c_n x^n = \sum_{n=0}^{\infty} c_n x^{n+r}$$

is a solution?

Method of Frobenius

□ If $x = x_0$ is a regular singular point of the differential equation $a_2(x)y'' + a_1(x)y' + a_0(x)y = 0$, then there exists at least one solution of the form

$$y = (x - x_0)^r \sum_{n=0}^{\infty} c_n (x - x_0)^n = \sum_{n=0}^{\infty} c_n (x - x_0)^{n+r},$$

where the number *r* is a constant to be determined. The series will converge on some interval of $0 < x - x_0 < R$.

Example:
$$3xy'' + y' - y = 0$$
 (1/3)

 \Box Solution: let $y = \sum_{n=0}^{\infty} c_n x^{n+r}$, we have

$$y' = \sum_{n=0}^{\infty} (n+r)c_n x^{n+r-1}$$
 and $y'' = \sum_{n=0}^{\infty} (n+r)(n+r-1)c_n x^{n+r-2}$.

Therefore,

$$3xy'' + y' - y$$

= $x^{r} \left[r(3r-2)c_{0}x^{-1} + \sum_{k=0}^{\infty} [(k+r+1)(3k+3r+1)c_{k+1} - c_{k}]x^{k} \right] = 0$
We have:
$$\begin{cases} r(3r-2)c_{0} = 0\\ (k+r+1)(3k+3r+1)c_{k+1} - c_{k} = 0, \ k = 0, 1, 2, \cdots \end{cases}$$

Example:
$$3xy'' + y' - y = 0$$
 (2/3)

□ Hence,

$$\rightarrow \begin{cases} r = 0, 2/3 \\ c_{k+1} = \frac{c_k}{(k+r+1)(3k+3r+1)}, k = 0, 1, 2, \cdots \end{cases}$$

Substituting r = 0 and r = 2/3 into the recurrence eq.,

$$\Rightarrow \begin{cases} r = 2/3, \ c_{k+1} = \frac{c_k}{(3(k+1)+2)(k+1)} \to c_n = \frac{c_0}{n! \, 5 \cdot 8 \cdot 11 \cdots (3n+2)} \\ r = 0, \quad c_{k+1} = \frac{c_k}{(k+1)(3(k+1)-2)} \to c_n = \frac{c_0}{n! \, 1 \cdot 4 \cdot 7 \cdots (3n-2)} \end{cases}$$

Example:
$$3xy'' + y' - y = 0$$
 (3/3)

 \Box Let $c_0 = 1$, we have two series solutions

$$\begin{bmatrix} y_1(x) = x^{2/3} \begin{bmatrix} 1 + \sum_{n=1}^{\infty} \frac{1}{n! 5 \cdot 8 \cdot 11 \cdots (3n+2)} x^n \end{bmatrix}, \\ y_2(x) = x^0 \begin{bmatrix} 1 + \sum_{n=1}^{\infty} \frac{1}{n! 1 \cdot 4 \cdot 7 \cdots (3n-2)} x^n \end{bmatrix}.$$

Since $y_1(x)$ and $y_2(x)$ are linearly independent on the entire axis, $y(x) = k_1y_1(x) + k_2y_2(x)$ is the general solution of the DE on any interval not containing the origin (note that 0^0 is undefined).

Indicial Equation

- The equation derived from the coefficient of the smallest degree of x in the Frobenius method is the indicial equation.
- □ The solutions of the indicial equation with respect to *r* are called the indicial roots.

Frobenius Series Solutions

Theorem: If x = 0 is a regular singular point of

$$x^{2}y'' + xp(x)y' + q(x)y = 0.$$

Let $\rho > 0$ denote the minimum of the radii of convergence of the power series

$$p(x) = \sum_{n=0}^{\infty} p_n x^n$$
 and $q(x) = \sum_{n=0}^{\infty} q_n x^n$

Let r_1 and r_2 be the (real) roots, with $r_1 \ge r_2$, of the indicial equation

$$r(r-1) + p_0 r + q_0 = 0.$$

Then, we have the following properties:

Frobenius Series Solutions

- (2/2)
- 1. For x > 0, there exists a solution of the form

$$y_1(x) = x^{r_1} \sum_{n=0}^{\infty} a_n x^n \quad (a_0 \neq 0)$$

corresponding to the larger root r_1 .

2. If $r_1 \neq r_2$ and $r_1 - r_2 \notin Z^+$, then there exists a 2nd linearly independent solution for x > 0 of the form

$$y_2(x) = x^{r_2} \sum_{n=0}^{\infty} b_n x^n \quad (b_0 \neq 0)$$

corresponding to the smaller root r_2 .

3. The radii of convergence of the solutions are at least ρ , the nearest distance to the nearby singular point.

Example:
$$2xy'' + (1 + x)y' + y = 0$$

□ Since $x^2y'' + \frac{1}{2} \cdot x(1+x) y' + \frac{1}{2} xy = 0$, $p(x) = \frac{1+x}{2}$ and $q(x) = \frac{x}{2} \rightarrow p_0 = \frac{1}{2}$ and $q_0 = 0 \rightarrow r^2 - \frac{r}{2} = 0$, $\rightarrow r = 0, \frac{1}{2}$.

For
$$r_1 = \frac{1}{2}$$
, let $y_1(x) = \sum_{n=0}^{\infty} a_n x^{n+r_1}$, then $a_n = \frac{(-1)^n a_0}{2^n n!}$

For
$$r_2 = 0$$
, let $y_2(x) = \sum_{n=0}^{\infty} b_n x^{n+r_2}$, we have
$$b_n = \frac{(-1)^n b_0}{1 \cdot 3 \cdot 5 \cdot 7 \cdots (2n-1)}.$$

The general solution is $y = c_1y_1(x) + c_2y_2(x)$.

Example:
$$xy'' + 2y' + xy = 0$$
 (1/3)

□ When $r_1 - r_2$ is a positive integer, the Frobenius solution is only guaranteed for r_1 . However, in this example, we still have two solutions even if $r_1 - r_2 = 1$.

The DE can be written as
$$y'' + \frac{2}{x}y' + \frac{x^2}{x^2}y = 0.$$

The indicial equation r(r-1) + 2r = 0 has roots 0, -1. Start with $r_2 = -1$, we have

$$y(x) = x^{-1} \sum_{n=0}^{\infty} c_n x^n = \sum_{n=0}^{\infty} c_n x^{n-1}$$

Hence,

$$\sum_{n=0}^{\infty} n(n-1)c_n x^{n-2} + \sum_{n=0}^{\infty} c_n x^n = 0. \rightarrow \text{Check the coefficients of } x^{-2} \text{ and } x^{-1}!$$

Example:
$$xy'' + 2y' + xy = 0$$
 (2/3)

The first two terms gives us $0 \cdot c_0 = 0$ and $0 \cdot c_1 = 0$, which means c_0 and c_1 can be arbitrary constants. Thus, the recurrence relation $c_n = -c_{n-2}/n(n-1)$, $n \ge 2$ can be divided into two groups of coefficients:

$$c_{2n} = \frac{(-1)^n c_0}{(2n)!}$$
 and $c_{2n+1} = \frac{(-1)^n c_1}{(2n+1)!}$, for $n \ge 1$.

Therefore, a general solution is

$$y(x) = x^{-1} \sum_{n=0}^{\infty} c_n x^n$$

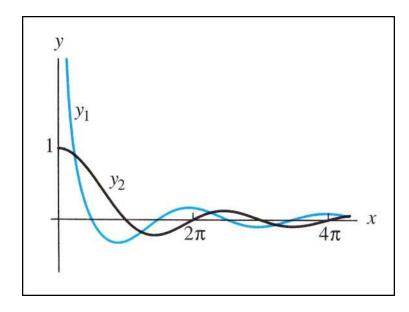
= $\frac{c_0}{x} \sum_{n=0}^{\infty} \frac{(-1)^n x^{2n}}{(2n)!} + \frac{c_1}{x} \sum_{n=0}^{\infty} \frac{(-1)^n x^{2n+1}}{(2n+1)!}$

Example:
$$xy'' + 2y' + xy = 0$$
 (3/3)

Now, if you pay attention, you will recognize that the solution is simply

$$y(x) = x^{-1}(c_0 \cos x + c_1 \sin x).$$

The graph of the solution is:



2nd Solution by Reduction of Order

□ If there is only one solution in Frobenius form for

y'' + P(x)y' + Q(x)y = 0,

we can find the 2nd solution by reduction of order.

Recall that the reduction of order formula tells us

$$y_2(x) = y_1(x) \int \frac{e^{-\int P(x)dx}}{y_1^2(x)} dx.$$

Summary of Indicial Roots (1/2)

□ Case I: r_1 and r_2 are distinct, $r_1 - r_2 \neq N$, for some integer $N \rightarrow$ exists two linearly independent solutions of the form

$$y_1(x) = \sum_{n=0}^{\infty} c_n x^{n+r_1}$$
 and $y_2(x) = \sum_{n=0}^{\infty} c_n x^{n+r_2}$.

□ Case II: $r_1 - r_2 = N$, for some integer $N \rightarrow$ exist two linearly independent solutions of the form

$$\begin{cases} y_1(x) = \sum_{n=0}^{\infty} c_n x^{n+r_1}, & c_0 \neq 0 \\ y_2(x) = C y_1(x) \ln x + \sum_{n=0}^{\infty} b_n x^{n+r_2}, & b_0 \neq 0 \end{cases}$$

Note that *C* could be zero.

Summary of Indicial Roots (2/2)

□ Case III: If $r_1 = r_2$, there exists two linearly independent solutions of the form

$$y_1(x) = \sum_{n=0}^{\infty} c_n x^{n+r_1}, \quad c_0 \neq 0$$
$$y_2(x) = y_1(x) \ln x + \sum_{n=1}^{\infty} b_n x^{n+r_1}$$

Bessel's Equations

 \Box Bessel's equation of order $v \ge 0$ is defined as

$$x^2y'' + xy' + (x^2 - v^2)y = 0.$$

The solutions are called Bessel functions of order v.

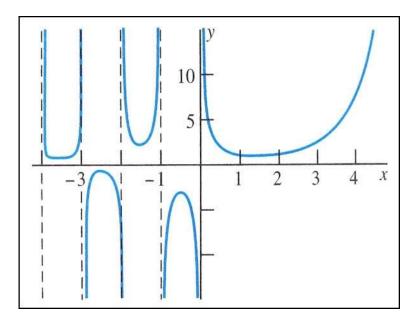
Bessel's functions first appear in 1764 when Euler was studying the vibration of drum membrane. Later, the functions appears in many physics problems, from fluid equations to planet motions.

Gamma Function $\Gamma(x)$

The gamma function (or generalized factorial function) is defined as

$$\Gamma(x) = \int_0^\infty t^{x-1} e^{-t} dt.$$

For x > 0, we have $\Gamma(x+1) = x \Gamma(x)$.



Solution of Bessel's Equation (1/2)

 \Box Let the solution be $y = \sum_{n=0}^{\infty} c_n x^{n+r}$, we have

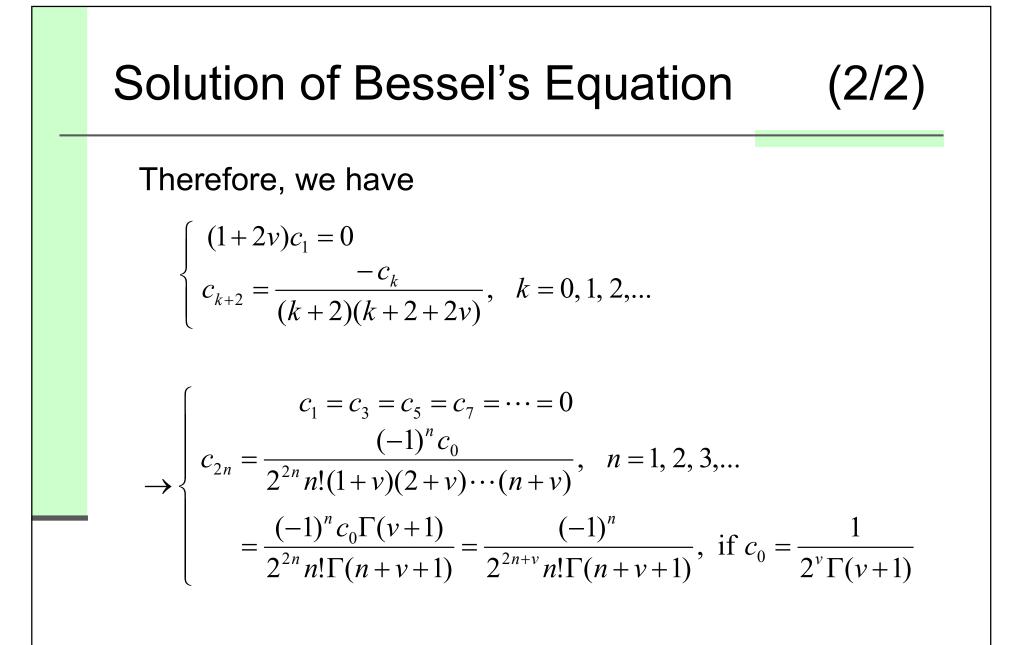
$$x^{2}y'' + xy' + (x^{2} - v^{2})y =$$

$$c_{0}(r^{2} - v^{2})x^{r} + x^{r}\sum_{n=1}^{\infty}c_{n}\left[(n+r)^{2} - v^{2}\right]x^{n} + x^{r}\sum_{n=0}^{\infty}c_{n}x^{n+2}$$

The indicial equation is $r^2 - v^2 = 0$, pick r = v

$$x^{\nu} \sum_{n=1}^{\infty} c_n n(n+2\nu) x^n + x^{\nu} \sum_{n=0}^{\infty} c_n x^{n+2}$$

$$= x^{\nu} \left[(1+2\nu)c_1 x + \sum_{k=0}^{\infty} \left[(k+2)(k+2+2\nu)c_{k+2} + c_k \right] x^{k+2} \right] = 0.$$



Bessel Functions of the 1st Kind (1/2)

□ The solutions of Bessel's Equation can be written as

$$J_{\nu}(x) = \sum_{n=0}^{\infty} c_{2n} x^{2n+\nu} = \sum_{n=0}^{\infty} \frac{(-1)^n}{n! \Gamma(1+\nu+n)} \left(\frac{x}{2}\right)^{2n+\nu}$$

Similarly, starting from r = -v, we have

$$J_{-\nu}(x) = \sum_{n=0}^{\infty} c'_{2n} x^{2n-\nu} = \sum_{n=0}^{\infty} \frac{(-1)^n}{n! \Gamma(1-\nu+n)} \left(\frac{x}{2}\right)^{2n-\nu}$$

 $J_{v}(x)$ and $J_{-v}(x)$ are called Bessel's functions of the first kind of order *v* and -v.

Bessel Functions of the 1st Kind (2/2)

□ Now, we want to find the general solution of the Bessel's DE. Notice that $r_1 - r_2 = 2v$:

- 1. If $2v \neq$ integer, then $J_{\nu}(x)$ and $J_{-\nu}(x)$ are linearly independent.
- 2. If 2v = 2m + 1, *m* is an integer, then $J_{m+1/2}(x)$ and $J_{-m-1/2}(x)$ are still linearly independent.
- 3. If 2v = 2m, *m* is an integer, then $J_m(x)$ and $J_{-m}(x)$ are linear dependent solutions of Bessel's DE. \rightarrow must find another solution!

$J_m \& J_{-m}$ are Linearly Dependent (1/2)

□ Proof:

Assume that v = m is an integer, we want to show that $J_{-m}(x) = (-1)^m J_m(x)$.

1) Perform change of index on $J_{-m}(x)$: $J_{-m}(x) = \sum_{n=0}^{\infty} \frac{(-1)^n}{n!\Gamma(1-m+n)} \left(\frac{x}{2}\right)^{2n-m}$

Let $2k+m = 2n-m \rightarrow k = n-m$ and n = k+m, we have

$$J_{-m}(x) = \sum_{k=-m}^{\infty} \frac{(-1)^{k+m}}{(k+m)!\Gamma(1+k)} \left(\frac{x}{2}\right)^{2k+m}$$

$J_m \& J_{-m}$ are Linearly Dependent (2/2)

2) Since $|\Gamma(x)| = \infty$, for x = 0, -1, -2, ..., we have

$$J_{-m}(x) = \sum_{k=0}^{\infty} \frac{(-1)^{k+m}}{(k+m)!\Gamma(1+k)} \left(\frac{x}{2}\right)^{2k+m}$$

3) Finally, note that

$$(k+m)!\Gamma(1+k) = [(k+m)(k+m-1) \dots (k+2)(k+1)] k! \Gamma(1+k)$$

= k! [(k+m)(k+m-1) \ldots (k+2)] \Gamma(2+k)
= k! \Gamma(1+m+k).

Therefore,

$$J_{-m}(x) = (-1)^m \sum_{k=0}^{\infty} \frac{(-1)^k}{k! \Gamma(1+m+k)} \left(\frac{x}{2}\right)^{2k+m} = (-1)^m J_m(x)$$

#

Bessel Functions of the 2nd Kind

□ If *v* is any non-integer number, we can apply linear combinations of $J_v(x)$ and $J_{-v}(x)$ to obtain another solution: $\int_{2}^{def} y_2(x) = Y_v(x) = \frac{\cos v \pi J_v(x) - J_{-v}(x)}{\sin v \pi}.$

For $m \in$ integer, $Y_m(x) = \lim_{v \to m} Y_v(x)$ still converges.

For any non-integer value of v, the general solution of Bessel's DE can also be written as

 $y = c_1 J_v(x) + c_2 Y_v(x).$

 $Y_{v}(x)$ is called the Bessel function of the 2nd kind.

Example: The Aging Spring

□ The DE for the free undamped motion of a mass on an aging spring is given by: $mx'' + ke^{-\alpha t}x = 0$. The change of variable,

$$s=\frac{2}{\alpha}\sqrt{\frac{k}{m}} e^{-\alpha t/2},$$

turns the DE into

$$s^2 \frac{d^2 x}{ds^2} + s \frac{dx}{ds} + s^2 x = 0.$$

Therefore, it's the Bessel DE with v = 0. The general solution is

$$x(t) = c_1 J_0 \left(\frac{2}{\alpha} \sqrt{\frac{k}{m}} e^{-\alpha t/2}\right) + c_2 Y_0 \left(\frac{2}{\alpha} \sqrt{\frac{k}{m}} e^{-\alpha t/2}\right).$$

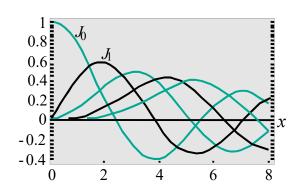
Properties of Bessel Functions

□ For m = 0, 1, 2, ..., we have:

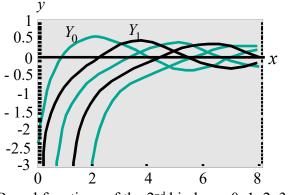
$$J_{-m}(x) = (-1)^m J_m(x)$$

- $\bullet J_m(-x) = (-1)^m J_m(x)$
- $J_m(0) = 0$ if m > 0; $J_m(0) = 1$, if m = 0

$$Iim_{x \to 0^+} Y_m(x) = -\infty$$



Bessel functions of the first kind, n = 0, 1, 2, 3, 4



Bessel functions of the 2^{nd} kind, n = 0, 1, 2, 3, 4

Bessel Functions with v = 0

□ When v = 0, we have $J_v(x) = J_{-v}(x)$, the 2nd solution can be obtained by Case III of the method of Frobenius: $y_1(x) = J_v(x)$, and

$$y_{2}(x) = \left(\sum_{n=0}^{\infty} \frac{(-1)^{n}}{(n!)^{2}} \left(\frac{x}{2}\right)^{2n}\right) \ln(x) + \sum_{n=1}^{\infty} b_{n} x^{n}.$$

Substitute $y_2(x)$ into the DE and solve for b_n , we have:

$$y_2(x) = \frac{2}{\pi} J_0(x) \left[\gamma + \ln \frac{x}{2} \right] - \frac{2}{\pi} \sum_{n=1}^{\infty} \frac{(-1)^n}{(n!)^2} \left(1 + \frac{1}{2} + \dots + \frac{1}{n} \right) \left(\frac{x}{2} \right)^{2n}$$

 $\gamma = 0.57721566$ is Euler's constant.

Differential Recurrence Relation

Bessel functions satisfy differential recurrence relations as follows:

•
$$xJ_{v}'(x) = vJ_{v}(x) - xJ_{v+1}(x)$$

•
$$xJ_{v'}(x) = xJ_{v-1}(x) - vJ_{v}(x)$$

□ To prove the relations, first, we have to show that

$$\frac{d}{dx} [x^{\nu} J_{\nu}(x)] = x^{\nu} J_{\nu-1}(x) \text{ and } \frac{d}{dx} [x^{-\nu} J_{\nu}(x)] = -x^{-\nu} J_{\nu+1}(x).$$

The recurrence relations can be derived easily, e.g.,

$$\frac{d}{dx} \Big[x^{-\nu} J_{\nu}(x) \Big] = -\nu x^{-\nu-1} J_{\nu}(x) + x^{-\nu} J_{\nu}'(x) = -x^{-\nu} J_{\nu+1}(x)$$

$$\rightarrow x J_{\nu}'(x) = \nu J_{\nu}(x) - x J_{\nu+1}(x).$$

Differentiation of $x^{\nu}J_{\nu}(x)$

□ Since

$$J_{\nu}(x) = \sum_{k=0}^{\infty} \frac{(-1)^{k}}{k! \Gamma(\nu+k+1)} \left(\frac{x}{2}\right)^{2k+\nu},$$

then

$$\frac{d}{dx} \left[x^{\nu} J_{\nu}(x) \right] = \frac{d}{dx} \sum_{k=0}^{\infty} \frac{(-1)^{k} x^{2k+2\nu}}{2^{2k+\nu} k! (\nu+k) \Gamma(\nu+k)}$$
$$= \sum_{k=0}^{\infty} \frac{(-1)^{k} x^{2k+2\nu-1}}{2^{2k+\nu-1} k! \Gamma(\nu+k)}$$
$$= x^{\nu} \sum_{k=0}^{\infty} \frac{(-1)^{k}}{k! \Gamma((\nu-1)+k+1)} \left(\frac{x}{2} \right)^{2k+(\nu-1)}$$
$$= x^{\nu} J_{\nu-1}(x)$$

Legendre's Equation

□ Legendre's equation of order α is the 2nd-order linear DE of the form

$$(1-x^2)y''-2xy'+\alpha(\alpha+1)y=0,$$

where the real number $\alpha > -1$. The only singular points of the Legendre's equation are at +1 and -1.

Solution of Legendre's Equation (1/2)

□ Since x = 0 is an ordinary point of the equation, substitute $y = \Sigma c_m x^m$ into the Legendre's equation, we have

$$c_{m+2} = -\frac{(\alpha - m)(\alpha + m + 1)}{(m+2)(m+1)}c_m, \ m \ge 0.$$

It can be shown that,

$$c_{2m} = (-1)^m \frac{\alpha(\alpha - 2)(\alpha - 4)\cdots(\alpha - 2m + 2)(\alpha + 1)(\alpha + 3)\cdots(\alpha + 2m - 1)}{(2m)!}c_0,$$

and

$$c_{2m+1} = (-1)^m \frac{(\alpha - 1)(\alpha - 3)\cdots(\alpha - 2m + 1)(\alpha + 2)(\alpha + 4)\cdots(\alpha + 2m)}{(2m+1)!}c_1.$$

Solution of Legendre's Equation (2/2)

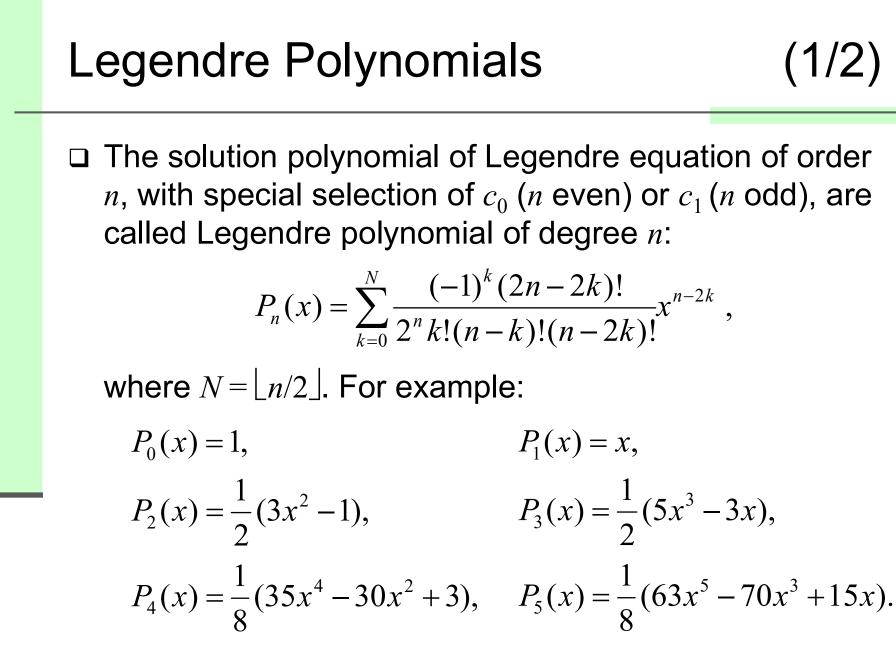
 \Box If $\alpha = n$, a non-negative integer, we have

$$y_1(x) = c_0 \left[1 - \frac{n(n+1)}{2!} x^2 + \frac{(n-2)n(n+1)(n+3)}{4!} x^4 - \frac{(n-4)(n-2)n(n+1)(n+3)(n+5)}{6!} x^6 + \cdots \right]$$

and

$$y_{2}(x) = c_{1} \left[x - \frac{(n-1)(n+2)}{3!} x^{3} + \frac{(n-3)(n-1)(n+2)(n+4)}{5!} x^{5} - \frac{(n-5)(n-3)(n-1)(n+2)(n+4)(n+6)}{7!} x^{7} + \cdots \right]$$

Notice that if *n* is an even integer, $y_1(x)$ terminates. When *n* is an odd integer, $y_2(x)$ terminates.



Legendre Polynomials

They are the solutions of

$$n = 0: \quad (1 - x^{2})y'' - 2xy' = 0$$

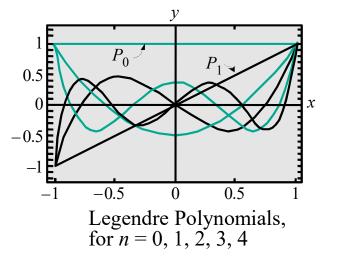
$$n = 1: \quad (1 - x^{2})y'' - 2xy' + 2y = 0$$

$$n = 2: \quad (1 - x^{2})y'' - 2xy' + 6y = 0$$

$$n = 3: \quad (1 - x^{2})y'' - 2xy' + 12y = 0$$

$$\vdots \qquad \vdots$$

Legendre polynomials are orthogonal over [-1, 1].



(2/2)